
ERRATA

Erratum: Joint probability distribution for Stark-broadening calculations
[Phys. Rev. A 34, 4091 (1986)]

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In plasmas, the emitting atoms and ions are under the influence of electric fields caused by electrons and ions [1,2]. The mean distance r_0 between ions is defined by [1,2]

$$(4\pi/3)r_0^3n=1 \quad (1)$$

and not

$$4\pi nr_0^3=1$$

as I have given in my paper. In the above equations, n stands for the ion density.

In Eq. (9), I should have ar_2/r_0 , and not ar_2 , because the argument of the exponential must be dimensionless.

Finally, the units for $Q(r,E)$ are $0.01 \epsilon_0^{-3}$ and not $0.01 (r_0\epsilon_0)^{-3}$. This is obvious from Eq. (4) by taking $A(\mathbf{r},\mathbf{E})=1$. However, it should be emphasized that the noted errors have no effect on my calculations since they were carried out with units such that $\epsilon=1$ and $r_0=1$.

I wish to acknowledge Ilunga S. K. Mwana Umbela for pointing out these errors.

[1] Hans. R. Griem, *Spectral Line Broadening by Plasmas* (Academic, New York, 1974).

[2] Hans R. Griem, *Plasma Spectroscopy* (McGraw-Hill Book Company, New York, 1964).

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Erratum: Dilute gas Couette flow: Theory and molecular-dynamics simulation
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Due to a mistake in a coefficient in one of our Maple three-dimensional routines, the expressions in the Appendix of the original article are wrong. The correct expressions are given below. Exact expressions are simpler when expressed in terms of $\gamma_s \equiv \frac{2}{5} \gamma$, where $\gamma = \tau v_x(z)'$ and $\tau = 5\sqrt{mT(z)}/\pi/(16\sigma^2 p)$. The hydrostatic pressure, like the adimensional shear rate γ , is uniform and, $p = P_{zz} - \frac{6}{5} \gamma P_{xz}$. Series expansions are in terms of γ itself. The expression $\Delta_3 = \sqrt{1 + 29\gamma_s^2 - 54\gamma_s^4}$ appears everywhere:

$$\begin{aligned} \frac{\eta}{\eta_0} &= \frac{2}{1 + 18\gamma_s^2 + \Delta_3} \\ &= 1 - \frac{13}{5} \gamma^2 + \frac{5282}{625} \gamma^4 + \dots, \end{aligned} \quad (1)$$

$$\begin{aligned}\frac{P_{xz}}{P_{zz}} &= \frac{-5\gamma_s}{1+3\gamma_s^2+\Delta_3} \\ &= -\gamma + \frac{7}{5}\gamma^3 - \frac{2282}{625}\gamma^5 + \dots,\end{aligned}\quad (2)$$

$$\begin{aligned}\frac{k_{xz}}{k_0} &= \frac{-\frac{35}{2}\gamma_s(1-\frac{9}{5}\gamma_s^2)}{1+\frac{15}{2}\gamma_s^2+(1-\frac{63}{4}\gamma_s^2)\Delta_3} \\ &= -\frac{7}{2}\gamma + \frac{1379}{500}\gamma^3 - \frac{87\,647}{5000}\gamma^5 + \dots,\end{aligned}\quad (3)$$

$$\begin{aligned}\frac{k_{zz}}{k_0} &= \frac{4}{1-54\gamma_s^2+3\Delta_3} \\ &= 1 + \frac{21}{50}\gamma^2 + \frac{6783}{2500}\gamma^4 + \dots,\end{aligned}\quad (4)$$

$$\begin{aligned}\frac{q_x}{\gamma q_z} &= -\frac{7-\frac{63}{5}\gamma_s^2}{1-3\gamma_s^2+\Delta_3} \\ &= -\frac{7}{2} + \frac{1057}{250}\gamma^2 - \frac{30\,653}{3125}\gamma^4 + \dots.\end{aligned}\quad (5)$$

The differential equation for the temperature profile $T(z)$ has the same form [Eq. (25)] as in the two-dimensional case except that instead of the constant K now appears K_3 ,

$$K_3 = \frac{\frac{64}{15}(8-63\gamma_s^2)(1+54\gamma_s^2)^2\gamma_s^2}{1+\frac{927}{14}\gamma_s^2+\frac{40\,689}{14}\gamma_s^4-3645\gamma_s^6+(1+\frac{891}{7}\gamma_s^2+486\gamma_s^4)\Delta_3} \pi\sigma^4 P_{zz}^2. \quad (6)$$

In the following expressions appear $f(\gamma)$, which we define so that $\sqrt{2K_3} = f(\gamma)\sigma^2 P_{zz}$.

When the differential equation for $T(z)$ is integrated the value T_{\max} at the center of the channel is introduced exactly as in Eq. (26) of the original paper, with K replaced by K_3 . The boundary and integral conditions render three conditions relating the constants P_{zz} , γ , and T_{\max} with the parameters externally imposed on the system like v_0 , T_0 , the number density N_{xy} of particles per unit area (in the Z direction), etc., and the parameters of the system like σ , L_z . The results are

$$\sqrt{T_{\max}-T_0} = \frac{5v_0}{32\gamma} \sqrt{\frac{m}{\pi}} \frac{f(\gamma)P_{zz}}{p}, \quad (7)$$

$$P_{zz} = \frac{2}{\sigma^2 L_z f(\gamma)} \left(T_0 \sqrt{\frac{T_{\max}-T_0}{T_0}} + T_{\max} \arctan \sqrt{\frac{T_{\max}-T_0}{T_0}} \right), \quad (8)$$

$$N_{xy} \sqrt{T_{\max}-T_0} = \frac{5v_0}{16\sigma^2\gamma} \sqrt{\frac{m}{\pi}} \arccos\left(\frac{2T_0-T_{\max}}{T_{\max}}\right). \quad (9)$$

We thank V. Garz3 for pointing out that there was something wrong with our Appendix.